

# Bragg Acceleration Structures

Amit Mizrahi and Levi Schächter



*Technion – Israel Institute of Technology  
Department of Electrical Engineering*

# *Outline*

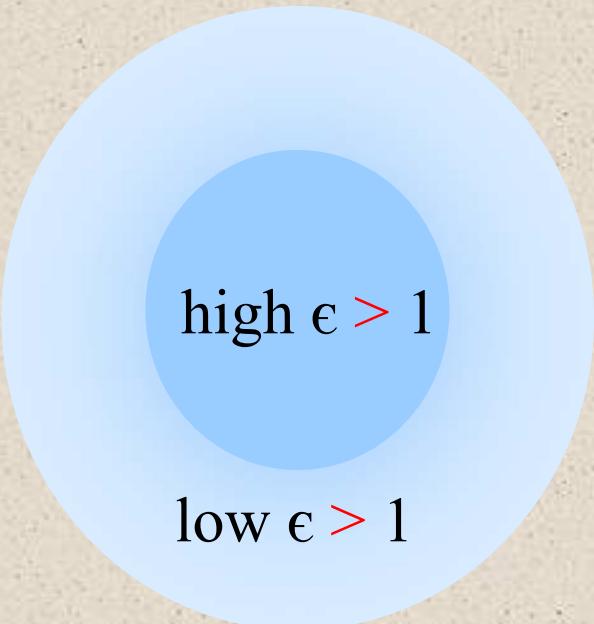
- Motivation and basic concept
- Confinement
- Accelerator parameters
- Wakes
- Efficiency

# Motivation

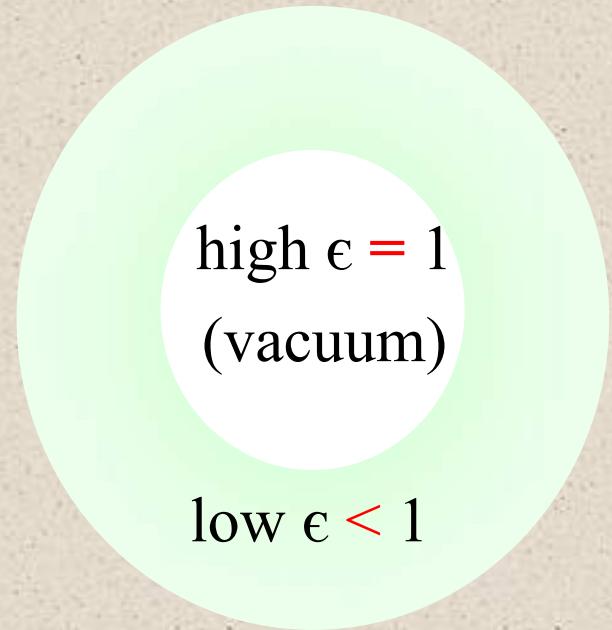
- Shorter and cheaper accelerators.
- Availability of high power lasers.
- Dielectrics sustain higher fields than metals.
- Fabrication: harness technology developed by communication or semiconductors industry.
- Need vacuum tunnel – confinement can not be achieved as in optical fibers – **Bragg waveguide!**

# *EM confinement*

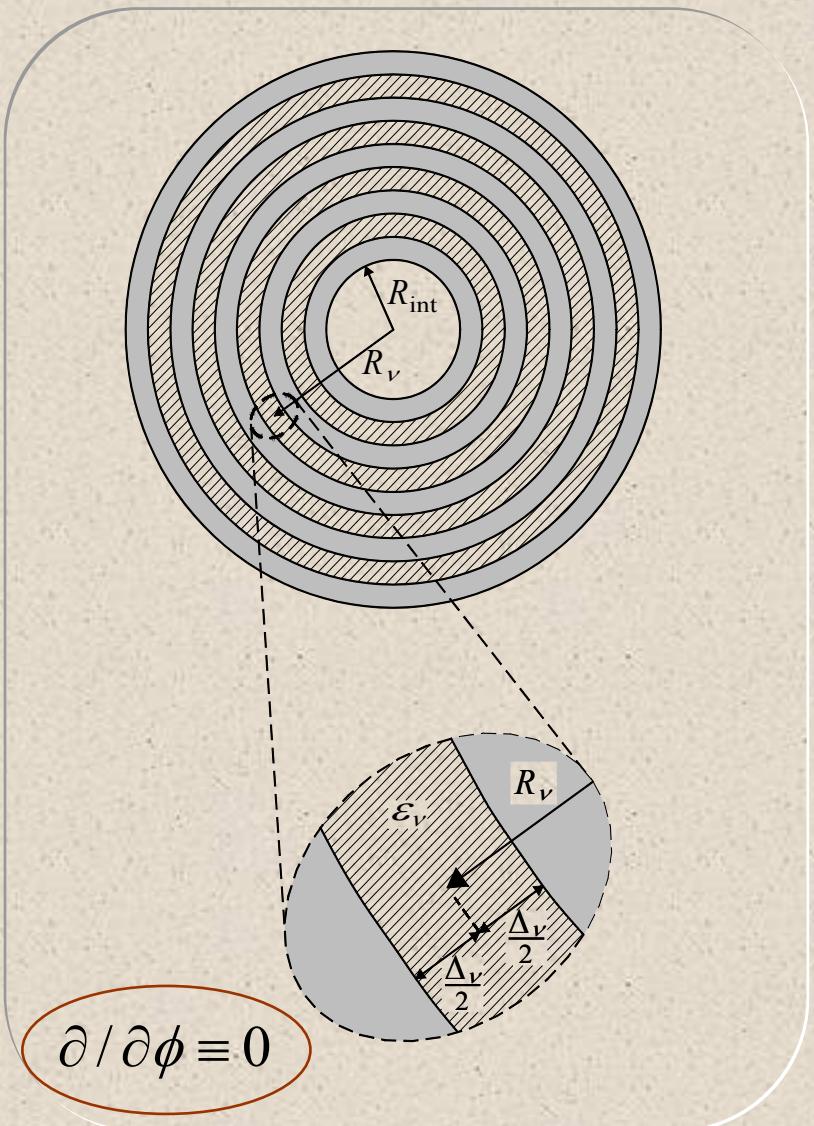
*Optical fiber*



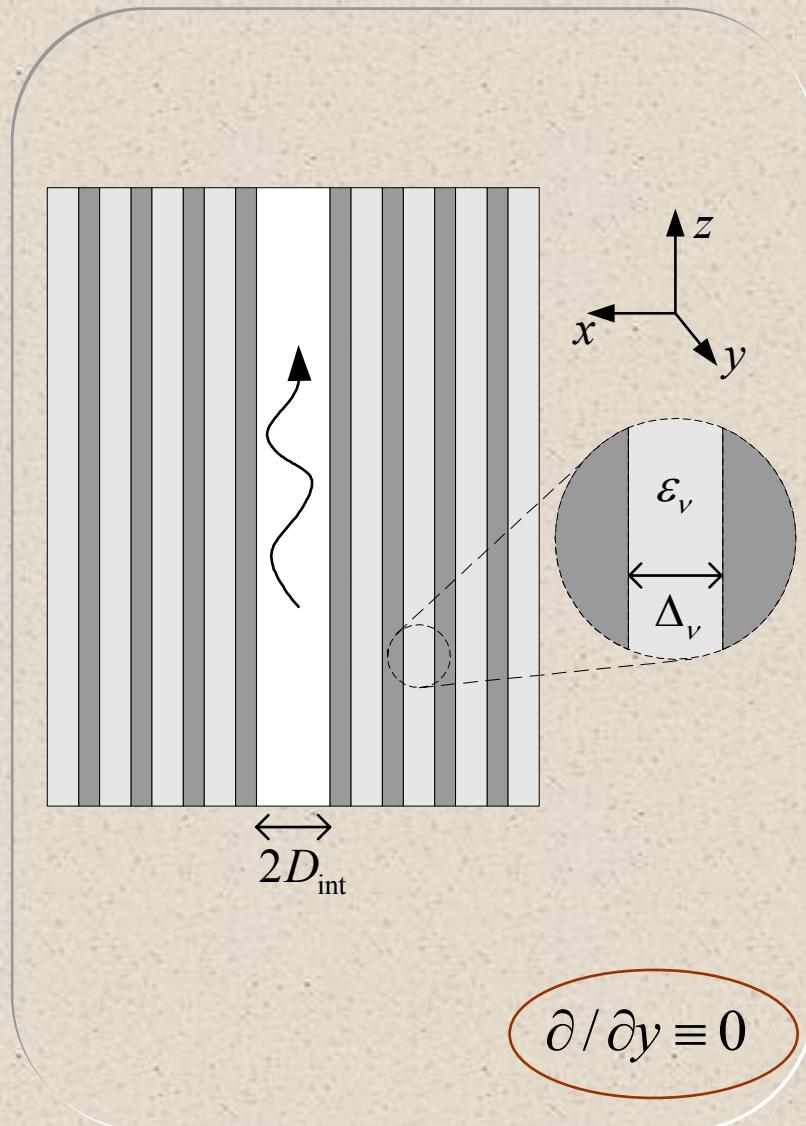
*Optical accelerator*



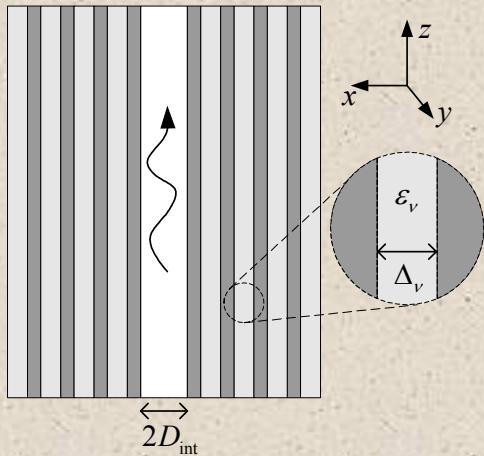
## Cylindrical Bragg Fiber



## Planar Bragg Waveguide



# *Evanescence wave: Infinite case*



T – Unit cell Transition matrix of incoming and outgoing amplitudes of transverse waves

Eigen-value problem  $|T - e^{-jKL} I| = 0$

$L$ –periodicity,  $K$ –propagation coefficient

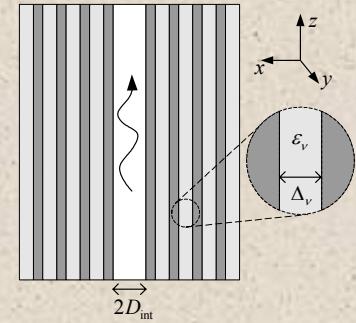
Dispersion relation  $\cos(KL) = \frac{1}{2}(T_{11} + T_{22})$

## Application to Bragg waveguides

- Yeh *et al.*, Opt. Commun. **19**, 427–430, (1976).
- Yeh *et al.*, JOSA, **68**, 1196–1201, (1978).

Confinement condition  $\left(\frac{T_{11} + T_{22}}{2}\right)^2 > 1$

# Optimal Confinement



Confinement condition

$$\left(\frac{T_{11} + T_{22}}{2}\right)^2 = \left(\frac{(Z_1 + Z_2)^2}{4Z_1 Z_2} \cos(\chi_1 + \chi_2) - \frac{(Z_1 - Z_2)^2}{4Z_1 Z_2} \cos(\chi_1 - \chi_2)\right)^2$$

$$\chi_{1,2} \triangleq 2\pi \frac{\Delta_{1,2}}{\lambda_0} \sqrt{\varepsilon_{1,2} - 1}$$



$$\Delta_{1,2} = \frac{\lambda_0}{4\sqrt{\varepsilon_{1,2} - 1}}$$

Quarter  $\lambda$  structure !!



$$|e^{-jKL}|^n = \begin{cases} \left(\frac{Z_1}{Z_2}\right)^n & Z_1 < Z_2 \\ \left(\frac{Z_2}{Z_1}\right)^n & Z_1 > Z_2 \end{cases}$$

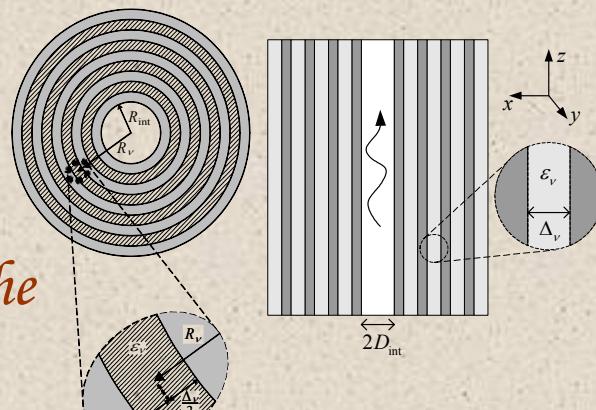


$$\left. \begin{array}{l} Z_1 > Z_2 \\ x \simeq nL \end{array} \right\} \Rightarrow \left(\frac{Z_2}{Z_1}\right)^{2n} \simeq \left(\frac{Z_2}{Z_1}\right)^{2x/L} \triangleq \exp\left(-2\frac{x}{x_c}\right)$$

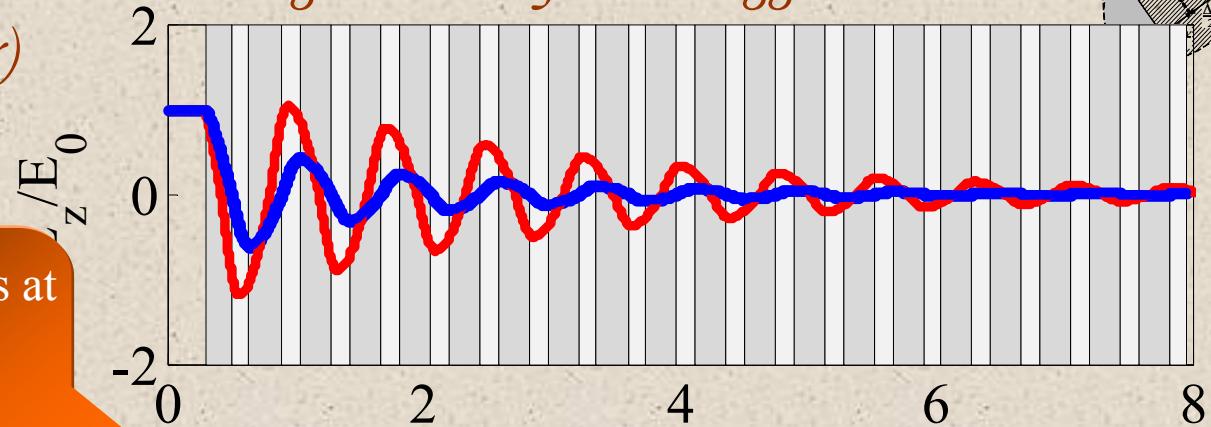
$$x_c = \frac{\lambda_0}{4} \left( \frac{1}{\sqrt{\varepsilon_1 - 1}} + \frac{1}{\sqrt{\varepsilon_2 - 1}} \right) \left| \ln^{-1} \left( \frac{\varepsilon_1 \sqrt{\varepsilon_2 - 1}}{\varepsilon_2 \sqrt{\varepsilon_1 - 1}} \right) \right|$$

# Field Confinement

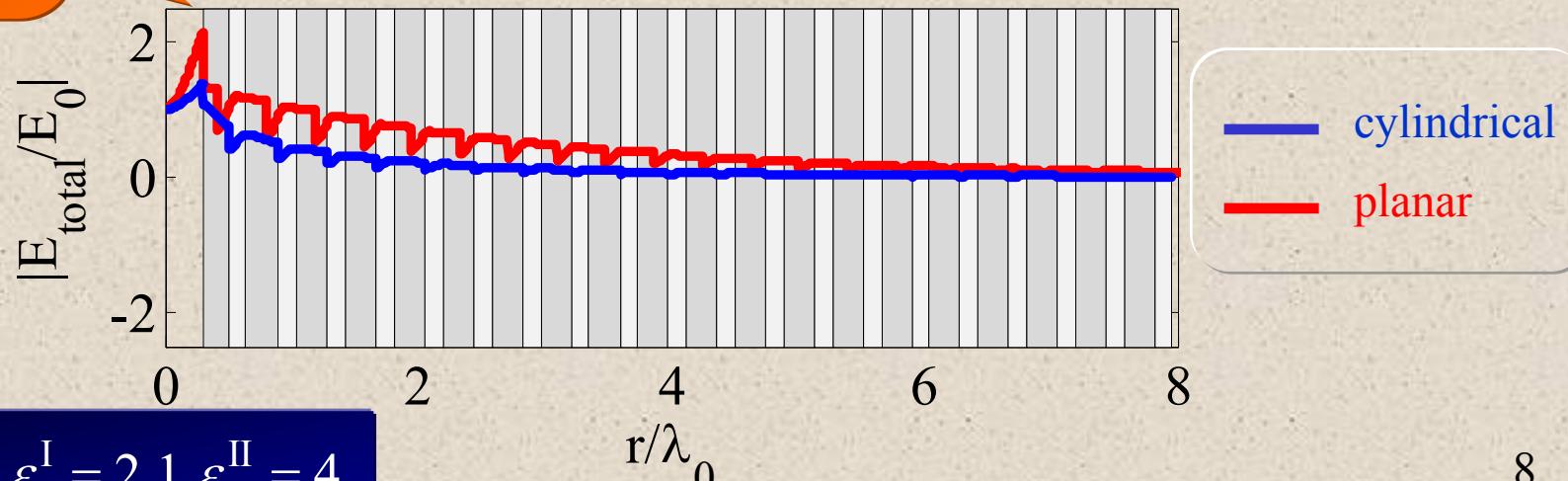
*First dielectric layer matches between the eigen-mode of the vacuum tunnel and the eigen-mode of the Bragg structure (stub tuner)*



Maximum is at  
vacuum-  
dielectric  
interface

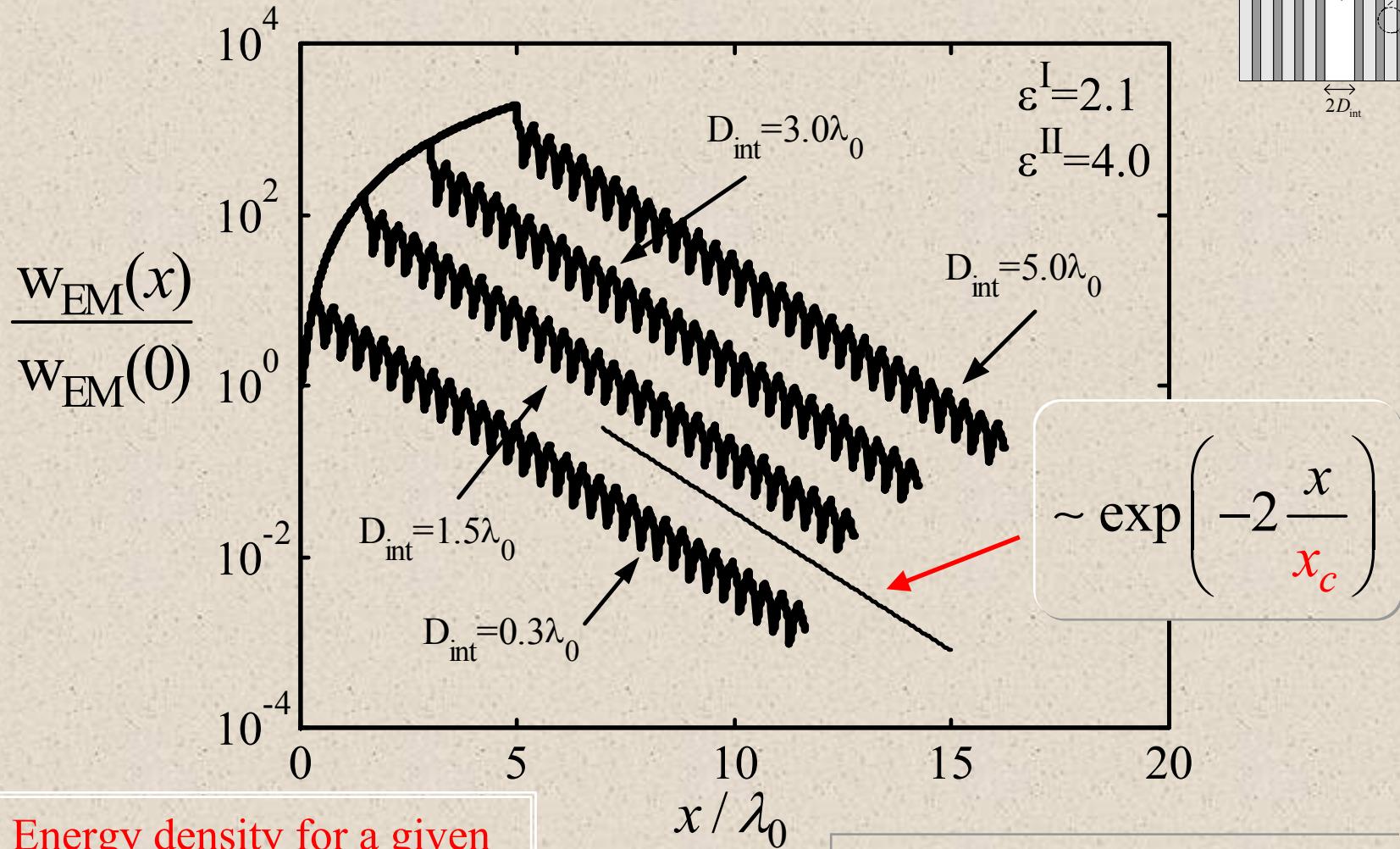
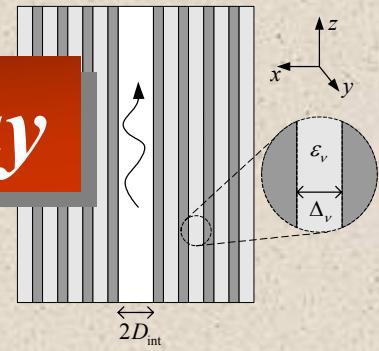


$E_z$  either peaks or diminishes at every discontinuity



$$D_{\text{int}} = 0.3\lambda_0, \epsilon^{\text{I}} = 2.1, \epsilon^{\text{II}} = 4$$

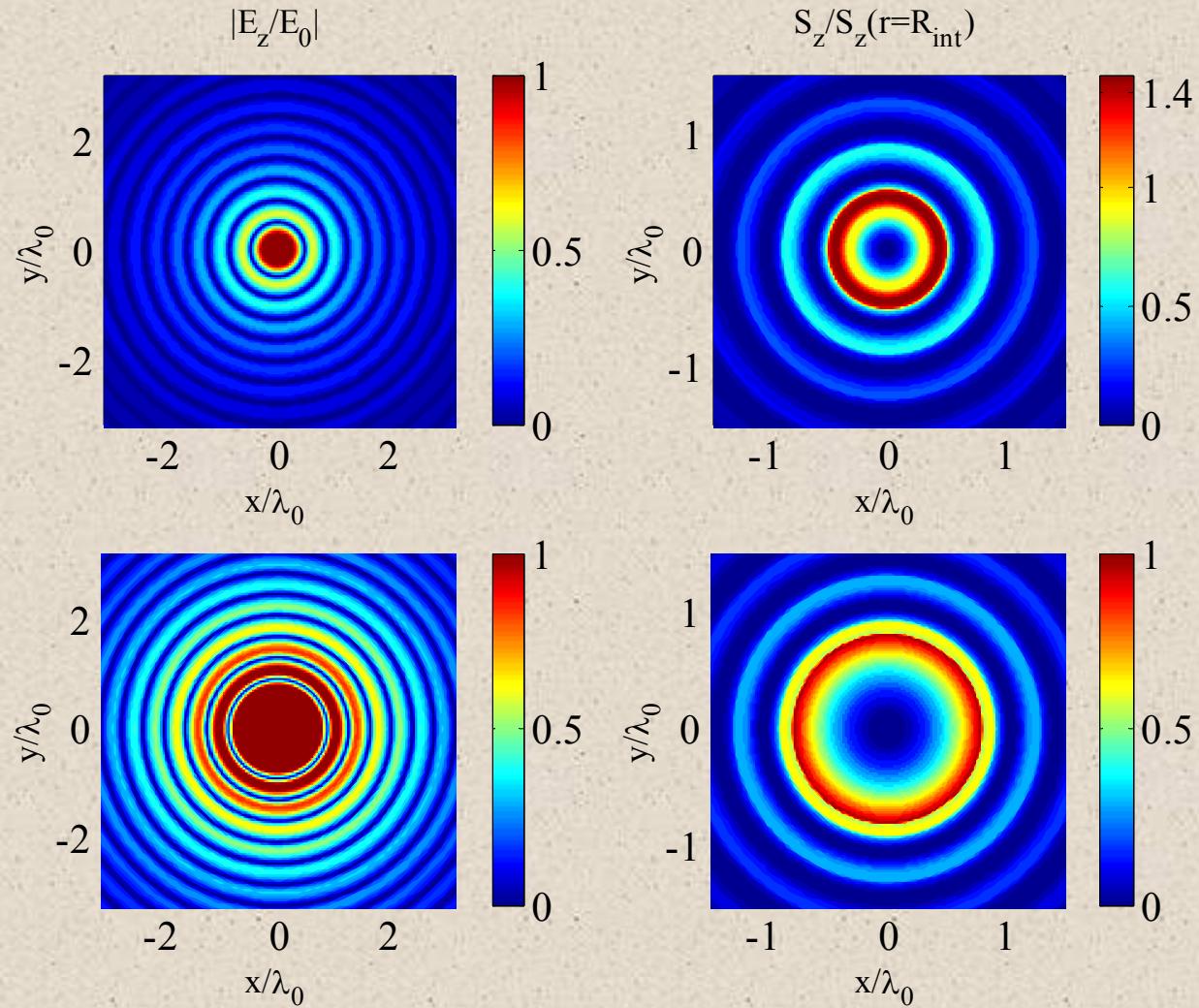
# Field Confinement – Energy Decay



Energy density for a given  $E_0$  increases for larger  $D_{int}/\lambda_0$  !! Breakdown.

$$x_c = \frac{\lambda_0}{4} \left( \frac{1}{\sqrt{\varepsilon_1 - 1}} + \frac{1}{\sqrt{\varepsilon_2 - 1}} \right) \left| \ln^{-1} \left( \frac{\varepsilon_1 \sqrt{\varepsilon_2 - 1}}{\varepsilon_2 \sqrt{\varepsilon_1 - 1}} \right) \right|$$

# Field Confinement – Cylindrical Case



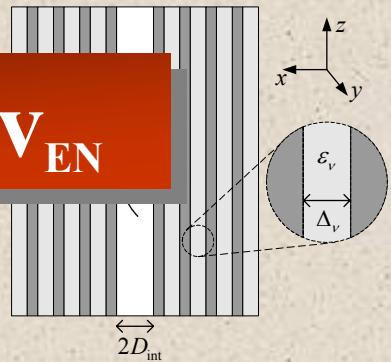
$$R_{\text{int}} = 0.3\lambda_0$$

$$\epsilon^I = 2.1$$

$$\epsilon^{II} = 4$$

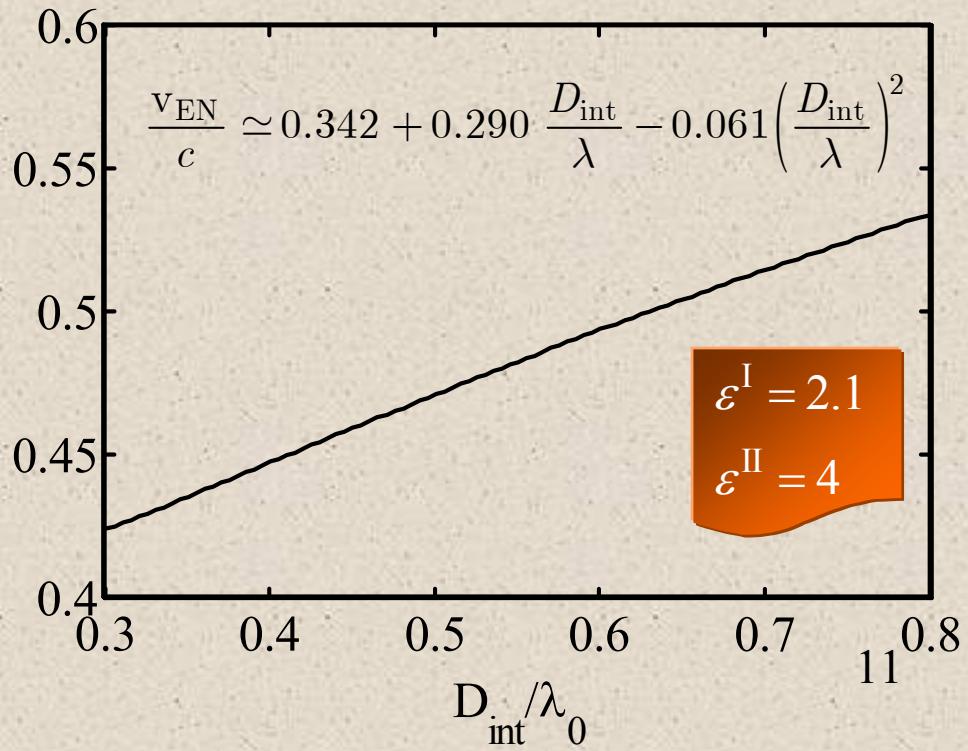
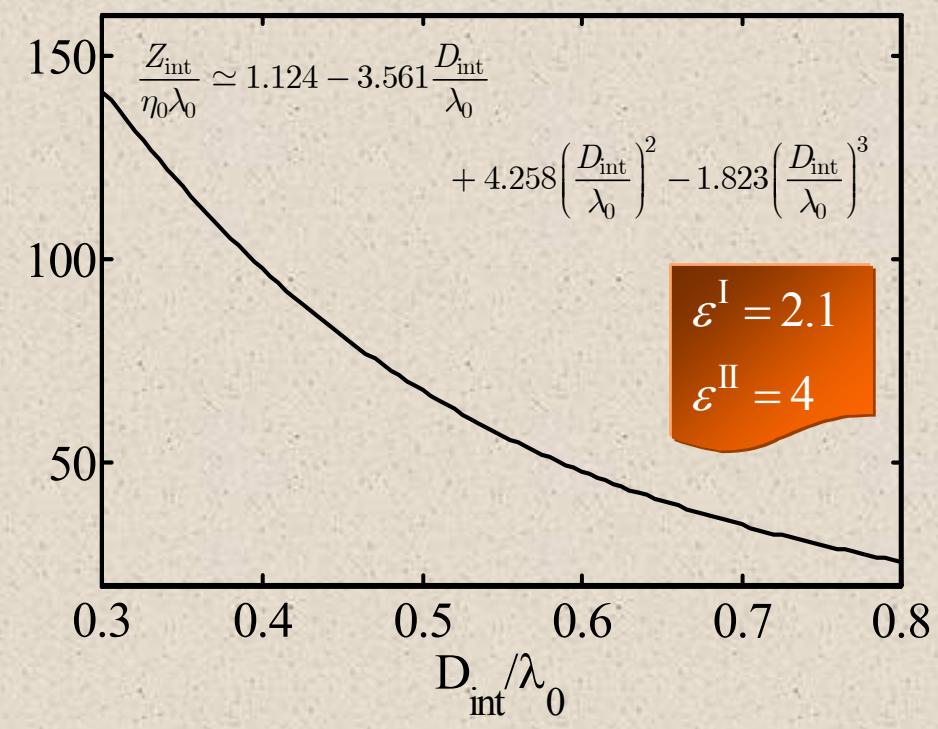
$$R_{\text{int}} = 0.8\lambda_0$$

# Accelerator Parameters – $Z_{\text{int}}$ and $v_{\text{EN}}$



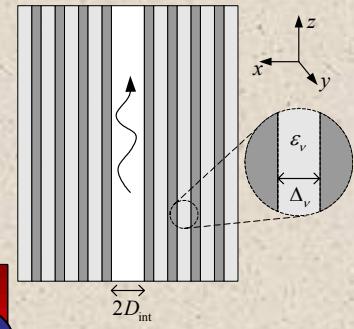
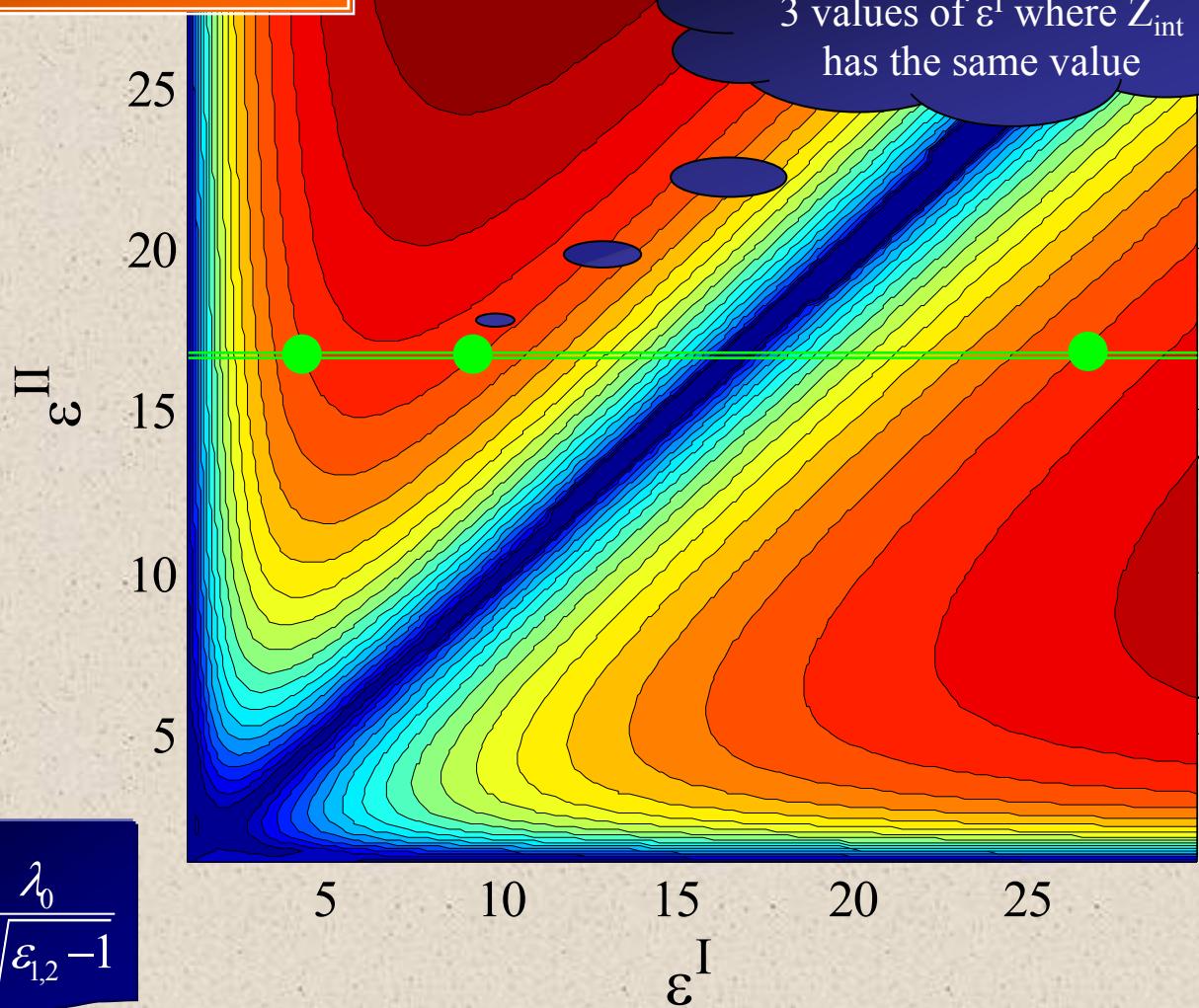
$$Z_{\text{int}} [\Omega m] \triangleq \frac{(\lambda_0 E_0)^2}{P}$$

$$\frac{v_{\text{EN}}}{c} \triangleq \frac{P}{cW}$$



# Accelerator Parameters – $Z_{\text{int}}$

High  $Z_{\text{int}}$  for large  $\varepsilon$



$$D_{\text{int}} = 0.3\lambda_0$$

$$\Delta_{1,2} = \frac{\lambda_0}{4\sqrt{\varepsilon_{1,2}-1}}$$

Non-symmetry: role of first layer

# Accelerator Parameters – $E_{\max}$

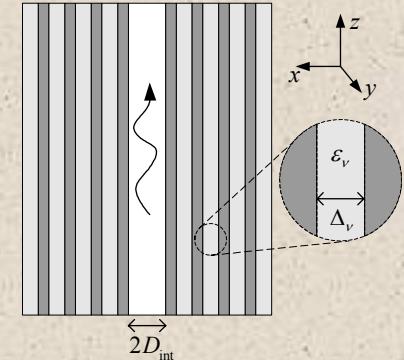
**Planar**

$$E_z = E_0 e^{-j \frac{\omega}{c} z}$$

$$E_x = j \frac{\omega}{c} x E_0 e^{-j \frac{\omega}{c} z}$$



$$\frac{E_{\max}}{E_0} = \sqrt{1 + \left( 2\pi \frac{D_{\text{int}}}{\lambda_0} \right)^2}$$



**Cylindrical**

$$E_z = E_0 e^{-j \frac{\omega}{c} z}$$

$$E_r = j \frac{\omega}{c} \frac{1}{2} r E_0 e^{-j \frac{\omega}{c} z}$$

$$1 \text{ [GV/m]}$$

$$2 \text{ [GV/m]}$$

$$\frac{E_{\text{acc}}}{E_{\max}} = 0.5$$

$$D_{\text{int}} \simeq 0.28\lambda_0$$

$$R_{\text{int}} \simeq 0.55\lambda_0$$



$$\frac{E_{\max}}{E_0} = \sqrt{1 + \left( \pi \frac{R_{\text{int}}}{\lambda_0} \right)^2}$$

# Wake-Field

Moving line charge

$q$  – charge per unit length

$$J_z(x, z, t) = -q v \delta(x) \delta(z - vt)$$

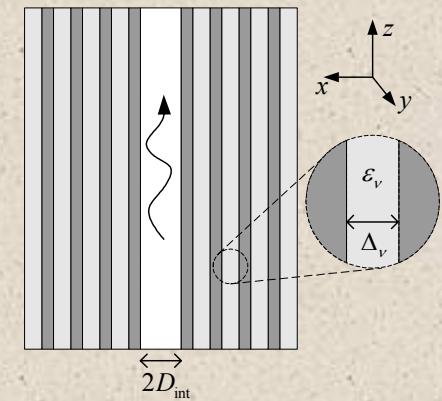
Primary potential

$$A_z^{(p)}(x, z, t) = -\frac{q \mu_0}{4\pi} \int_{-\infty}^{\infty} d\omega e^{j\omega(t-z/v)} \frac{1}{\Gamma} e^{-\Gamma|x|}$$

Secondary potential

$$A_z^{(s)}(x, z, t) = -\frac{q \mu_0}{4\pi} \int_{-\infty}^{\infty} d\omega e^{j\omega(t-z/v)} \frac{1}{\Gamma} \begin{cases} B_0 \cosh(\Gamma x) & |x| < D_{\text{int}} \\ -\frac{\varepsilon_1}{\gamma^2 \beta^2 (\varepsilon_1 - \beta^{-2})} (C_0 e^{-j\Lambda x} + D_0 e^{j\Lambda x}) & x > D_{\text{int}} \end{cases}$$

All waves have  $k_z = \omega/v$  !!



$$\beta \triangleq v/c$$

$$\gamma \triangleq 1/\sqrt{1-v^2/c^2}$$

$$\Gamma \triangleq \frac{|\omega|}{c\gamma\beta}$$

$$\Lambda \triangleq \frac{|\omega|}{c} \sqrt{\varepsilon_1 - \beta^{-2}}$$

$$R_{11} \triangleq \frac{D_0}{C_0} e^{2j\Lambda D_{\text{int}}}$$

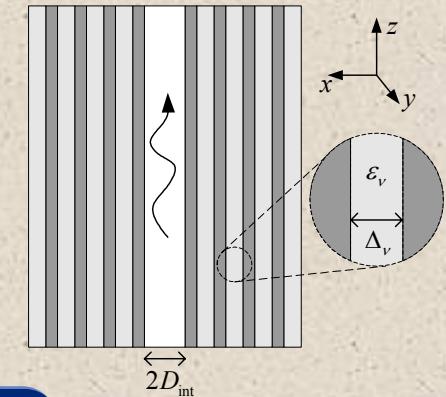
# Wake-Field On Axis

Ultra-relativistic wake-field on axis  $\gamma \rightarrow \infty$

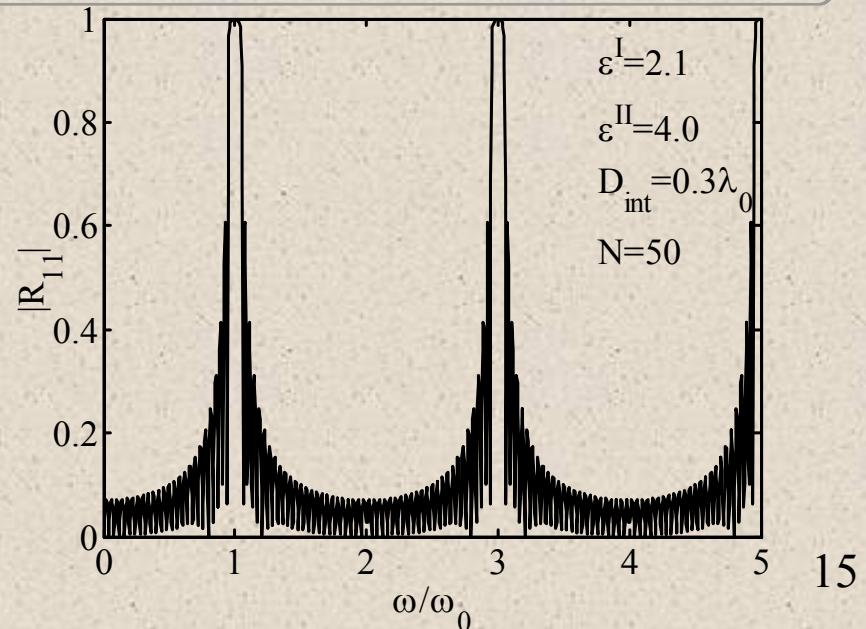
$$E_z^{(s)}(\bar{\tau}) = \frac{q}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{1}{2} \int_{-\infty}^{\infty} d\bar{\omega} e^{j\bar{\omega}\bar{\tau}} \frac{1 + R_{11}(\bar{\omega})}{(1 + j\bar{\omega}) - (1 - j\bar{\omega})R_{11}(\bar{\omega})}$$

$$\bar{\omega} \triangleq \frac{\omega}{c} D_{\text{int}} \frac{\sqrt{\epsilon_1 - 1}}{\epsilon_1}$$

$$\bar{\tau} \triangleq \left( t - \frac{z}{c} \right) \frac{c}{D_{\text{int}}} \frac{\epsilon_1}{\sqrt{\epsilon_1 - 1}}$$



Reflection coefficient (analytic expression)

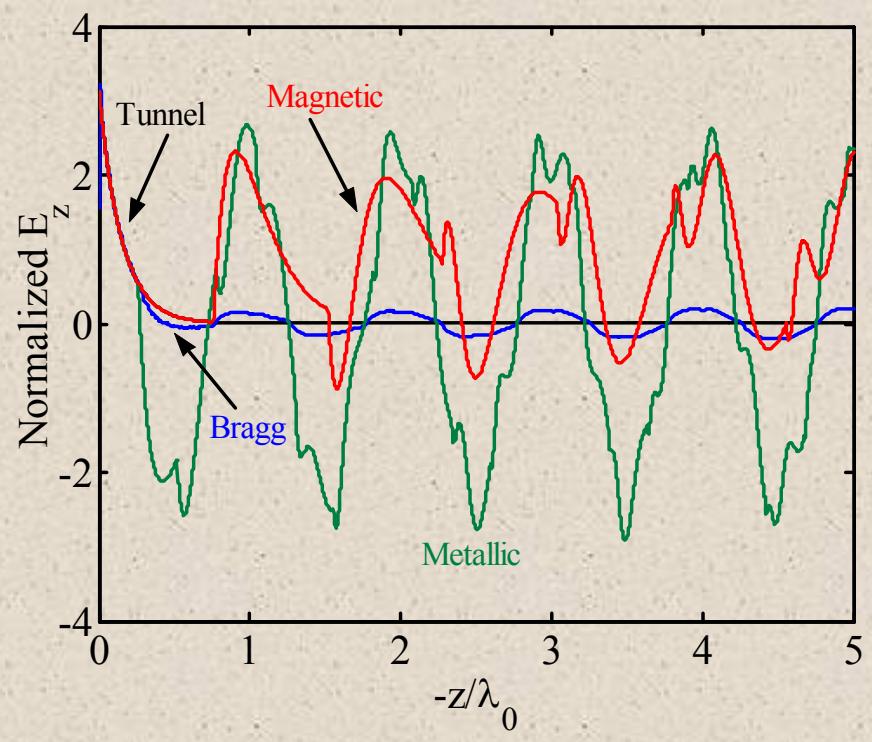
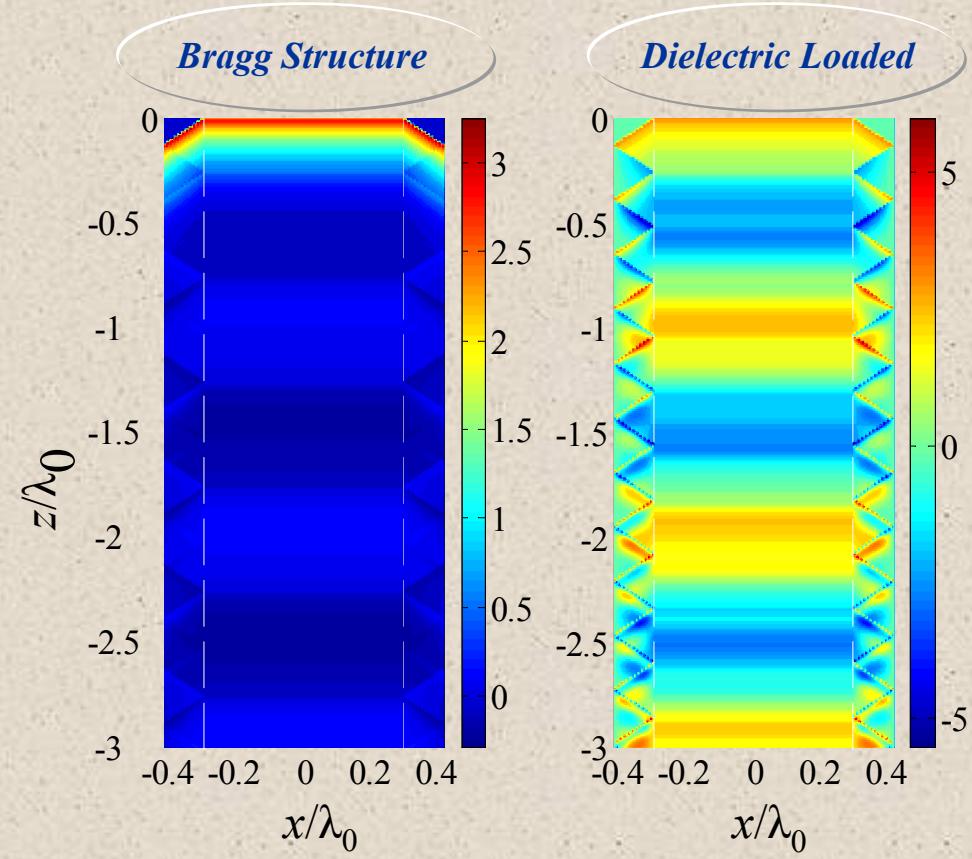
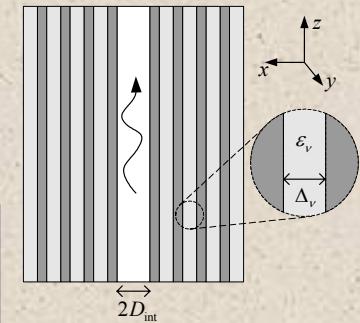


# Wake-field- $t=0$ picture

$$D_{\text{int}} = 0.3\lambda_0$$

$$\epsilon^{\text{I}} = 2.1$$

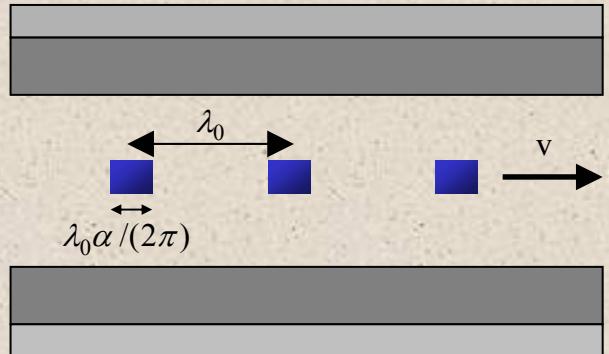
$$\epsilon^{\text{II}} = 4$$



# Emitted Power – Qualitative Approach

one line-charge

$$P = \frac{vq^2}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{\pi}{2}$$



Total power  
neglecting  
mutual effects.  
 $q = N_{el}q_{el}, \alpha=0$

$$P = \frac{vq_{el}^2 N_{el}^2 / M^2}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{\pi}{2} \times M = \frac{vq_{el}^2}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{N_{el}^2}{M} \times \frac{\pi}{2}$$

One micro-bunch  
(line-charge)

$$P = \frac{vq_{el}^2}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{\pi}{2} \times N_{el}^2$$

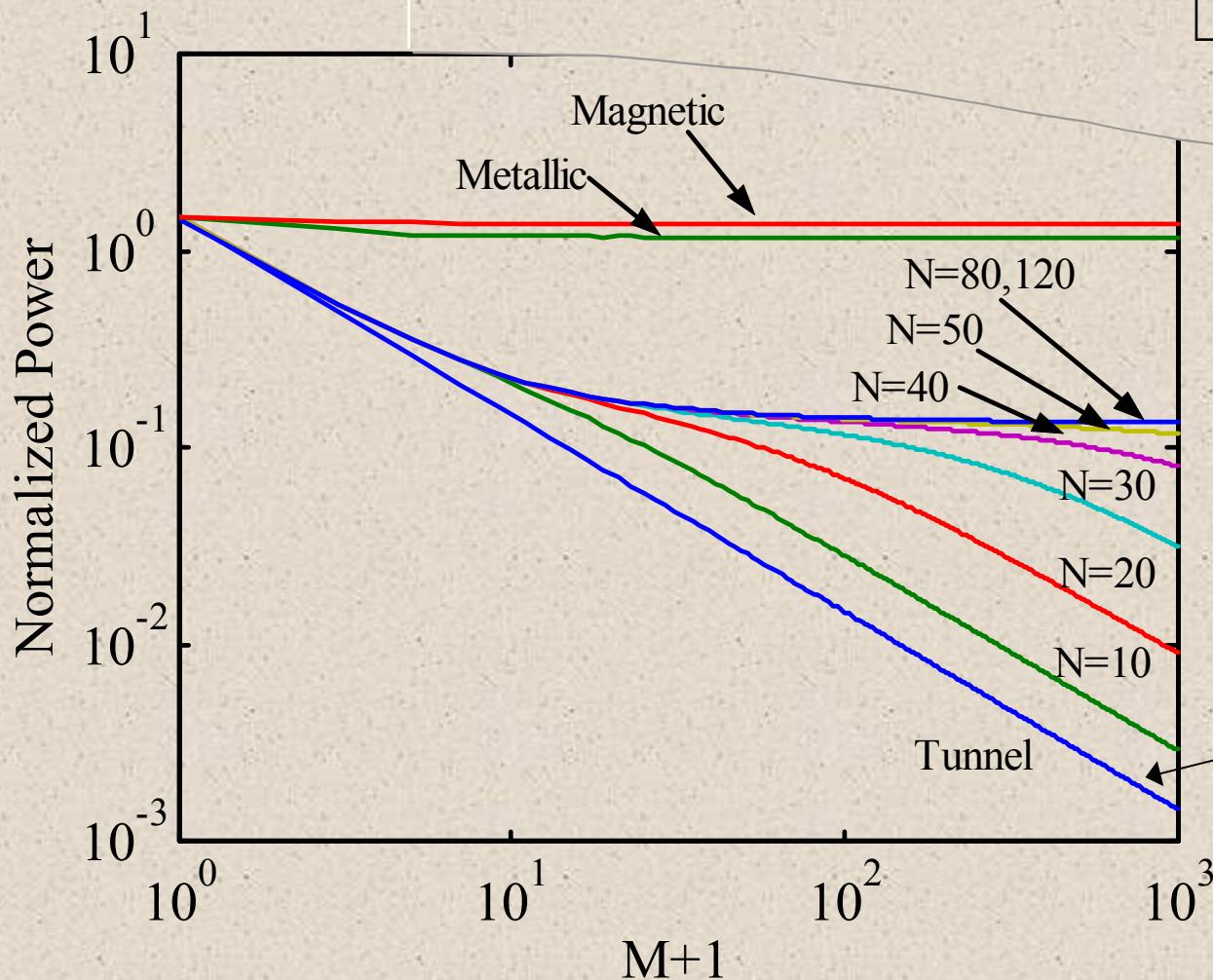
$\sim 1/M$

Randomly  
distributed

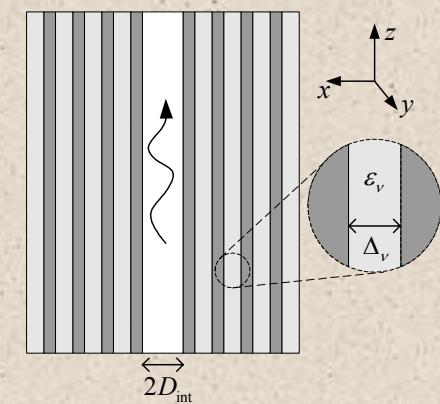
$$P = \frac{vq_{el}^2}{2\pi\epsilon_0 D_{\text{int}}} \times \frac{\pi}{2} \times N_{el}$$

$$P = \frac{vq^2}{2\pi\epsilon_0 D_{\text{int}}} \frac{1}{2} \int_{-\infty}^{\infty} d\bar{\omega} \frac{1 + R_{11}(\bar{\omega})}{(1 + j\bar{\omega}) - (1 - j\bar{\omega})R_{11}(\bar{\omega})}$$

## Emitted Power



$$\times \text{sinc}^2 \left[ \frac{\alpha}{2} \frac{\bar{\omega}}{\bar{\omega}_0} \right] \frac{\text{sinc}^2 \left[ \pi \frac{\bar{\omega}}{\bar{\omega}_0} (M+1) \right]}{\text{sinc}^2 \left[ \pi \frac{\bar{\omega}}{\bar{\omega}_0} \right]}$$



# Efficiency

- *Wake parameter:*

Decelerating field for a given charge .....  $E_{\text{dec}} \equiv \kappa q$

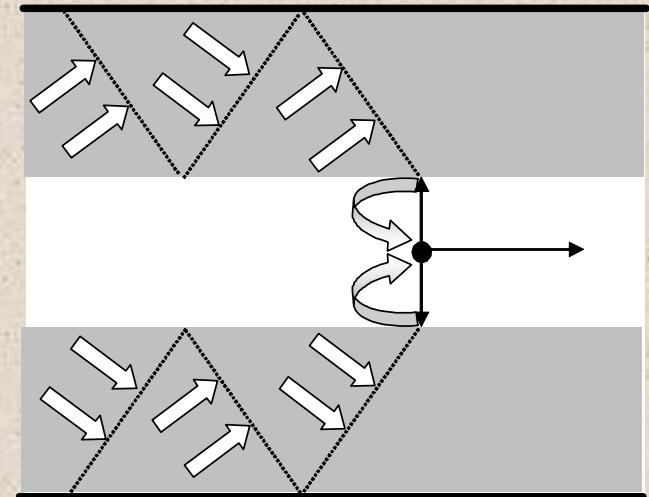
- *Beam-loading parameter:*

Beam-loading of the accelerating mode .....  $E_{\text{dec}}^{(\text{F})} \equiv \kappa_1 q$

$$E_z(r=0, \tau = t - z/c) \approx q\kappa \sum_{n=1}^{\infty} W_n \cos(\omega_n \tau) 2h(\tau)$$

$$W_n = \left[ \frac{2J_1(p_n R_b / R_{\text{ext}})}{p_n J_1(p_n)} \right]^2$$

$$\sum_{n=1}^{\infty} W_n = 1$$

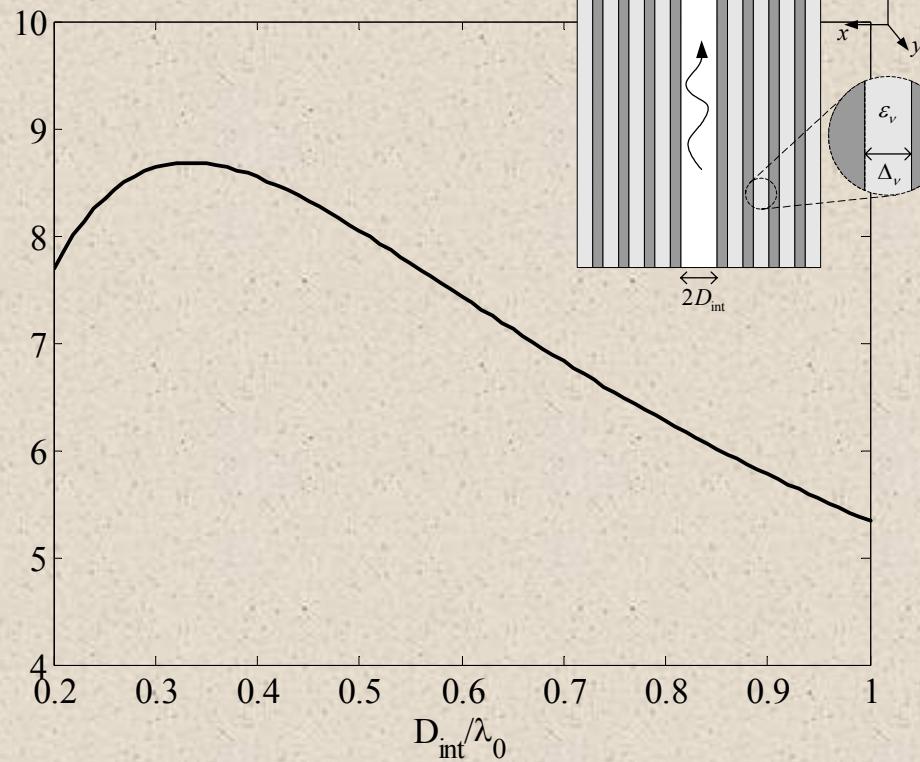
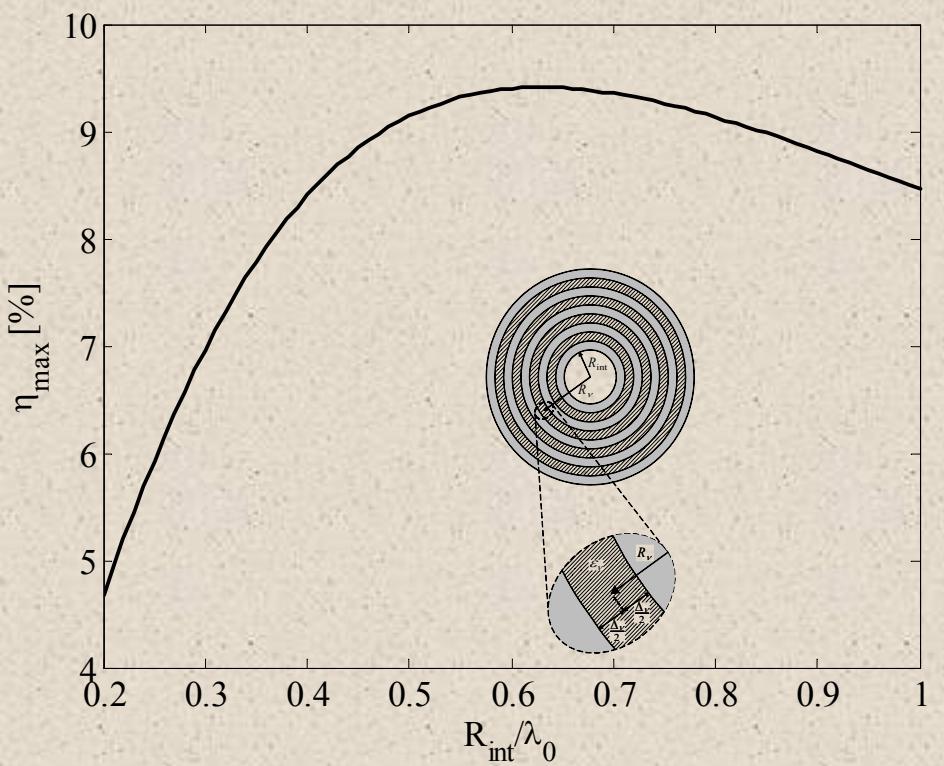


$$\kappa_1 \equiv \kappa W_1$$

$$\kappa_1 = \frac{\beta_{\text{gr}}}{1 - \beta_{\text{gr}}} \frac{Z_{\text{int}}}{\sqrt{\mu_0 / \epsilon_0}} \frac{\pi}{4\pi\epsilon_0\lambda^2}$$

$$\eta_{\max} = \frac{\kappa_1}{\kappa}$$

# Efficiency



$$\begin{aligned}\epsilon^{\text{I}} &= 2.1 \\ \epsilon^{\text{II}} &= 4\end{aligned}$$

# *Summary*

- Detailed design of Bragg acceleration structures – theoretical feasibility was introduced (PRE 2004)
- Structure parameters (interaction impedance, group velocity, maximal field).
- “Better” materials can dramatically improve performance.
- Analysis of wake-field – power decreases with the number of micro-bunches, and increases with the number of layers.
- Efficiency has an optimum within an internal dimension range of  $0.3 \div 0.8 \lambda_0$ .

# Summary - parameters

Silica-Zirconia structure,  $R_{\text{int}}, D_{\text{int}} = 0.3 \div 0.8 \lambda_0$

$$\varepsilon^{\text{I}} = 2.1$$

$$\varepsilon^{\text{II}} = 4$$

	<b>Planar</b>	<b>Cylindrical</b>
$Z_{\text{int}} [\Omega]$	$147 \div 25 \lambda_0 / \Delta_y$	$268 \div 37$
$\beta_{\text{gr}}$	$0.42 \div 0.53$	$0.41 \div 0.48$
$E_{\text{acc}}/E_{\text{max}}$	$0.47 \div 0.20$	$0.73 \div 0.37$
$\kappa$	$1/4 \varepsilon_0 D_{\text{int}}$	$1/2 \pi \varepsilon_0 R_{\text{int}}^2$
$\eta_{\text{max}} [\%]$	$8.7 \div 6.23$	$7 \div 9.4$

